



Danish Quantum Optics Center

Light-Atoms Quantum Interface for Continuous Variables

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QUICOV

Content

- Continuous variables for polarized light and oriented atoms
- Hamiltonian of light-atoms interaction
- Light-atoms interface via measurement induced back action: interface with and without teleportation
- Experiment with vacuum and squeezed vacuum input – recording a single quadrature

Experiment: recording in the ground state coherence ($\tau=1.5\text{msec}$)

Schori, Julsgaard, Sorensen, Polzik

Phys. Rev. Lett. **89**, 057903, July 29, 2002

quant-ph/0203023

Theory - interspecies teleportation

Kuzmich, Polzik

Phys. Rev. Lett. **85**, 5639 (2000)

Recording w/out teleportation

Kuzmich, Polzik to appear in

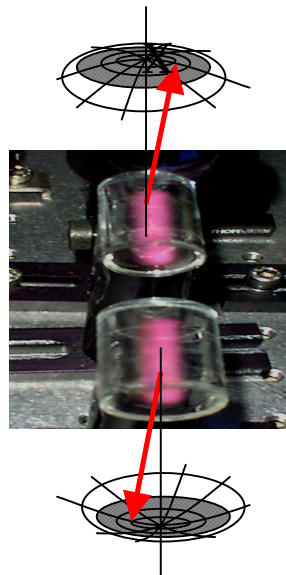
“QI with continuous variables”

ed. Braunstein, Pati. Kluwer

Related work

Julsgaard, Kozhekin, Polzik

Nature **413**, 400 (2001).

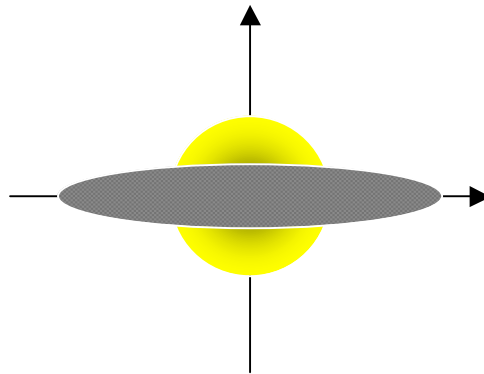


Continuous variables - light

Quadratures of light

$$X = \frac{1}{\sqrt{2}} (a^+ + a), P = \frac{i}{\sqrt{2}} (a^+ - a)$$

$$[\hat{X}, \hat{P}] = i$$

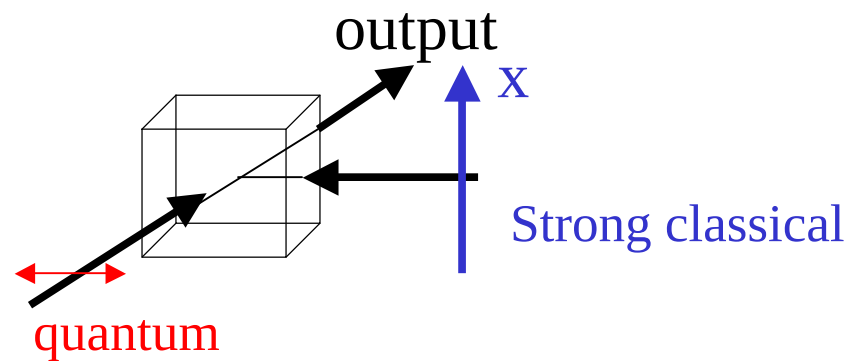


Can states of light like that be faithfully recorded in atoms?

Can such states be distinguished in the read out from atoms?

NB: the states are not orthogonal

Continuous variables – polarization states

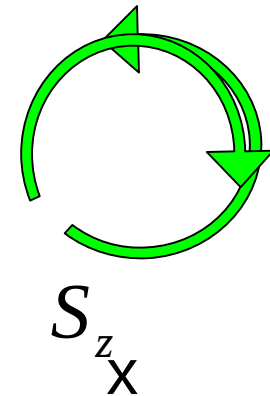


Output polarization –
Stokes parameters

$$[\hat{S}_z, \hat{S}_y] = iS_x$$

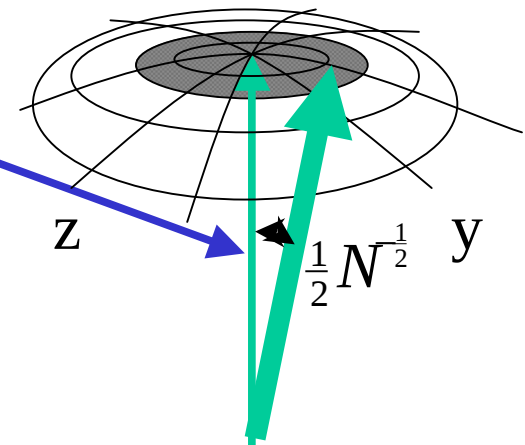
$$S_x = a_v^+ a_v - a_h^+ a_h$$

Diagram illustrating the Stokes parameters S_x and S_y . S_x is represented by a vertical blue arrow labeled v and a horizontal green arrow labeled h . S_y is represented by two green arrows at -45° and 45° .



Large = classical

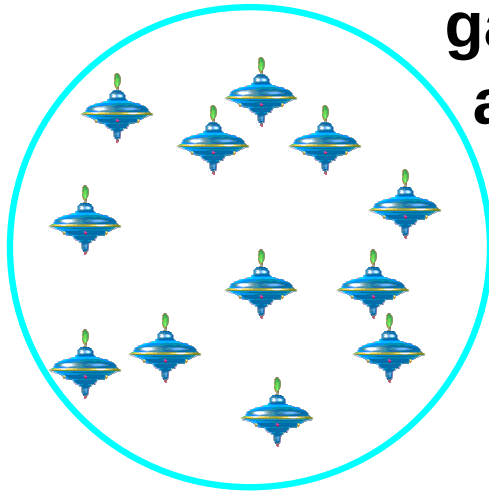
$$\hat{X} = \frac{\hat{S}_z}{\sqrt{S_x}}, P = \frac{\hat{S}_y}{\sqrt{S_x}}$$



Macroscopic spin ensemble –

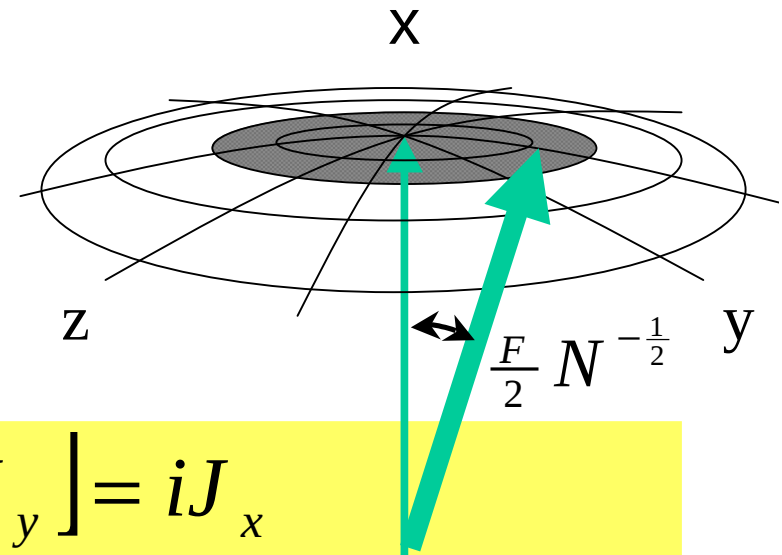
coherent spin state

gas sample
at room T



X

State
preparation
(optical
pumping)



$$[J_z, J_y] = iJ_x$$

$$\delta J_z^2 = \delta J_y^2 = \frac{1}{2} J_x = \frac{F}{2} N$$

Cesium $6P_{3/2}$

$$J_x = N \sum_{m=(-4, \dots, 4)} m \rho_{mm}$$

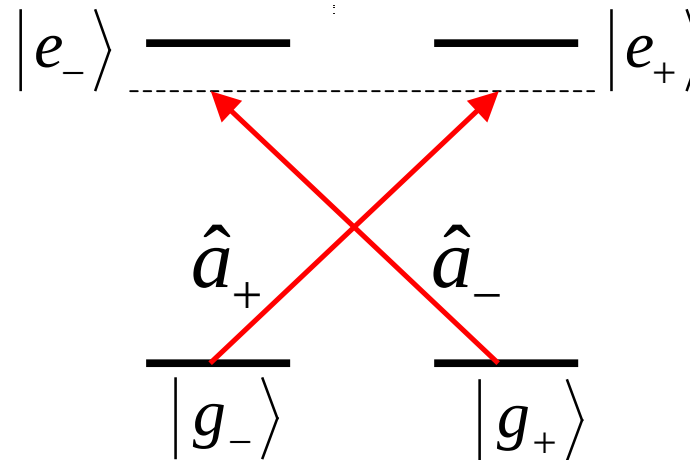
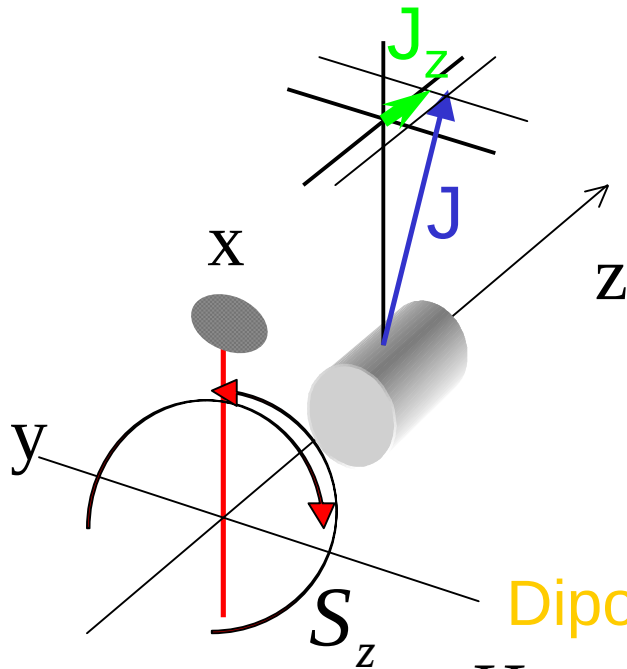
$6S_{1/2}, F=4$

Coherences determine J_z, J_y

Light / Atom - Interaction

Hamiltonian:

$$H = a S_z J_z$$



Dipole interaction Hamiltonian:

$$H \propto dE \propto a_+ |g_-\rangle \langle e_+| + a_- |g_+\rangle \langle e_-| + h.c.$$

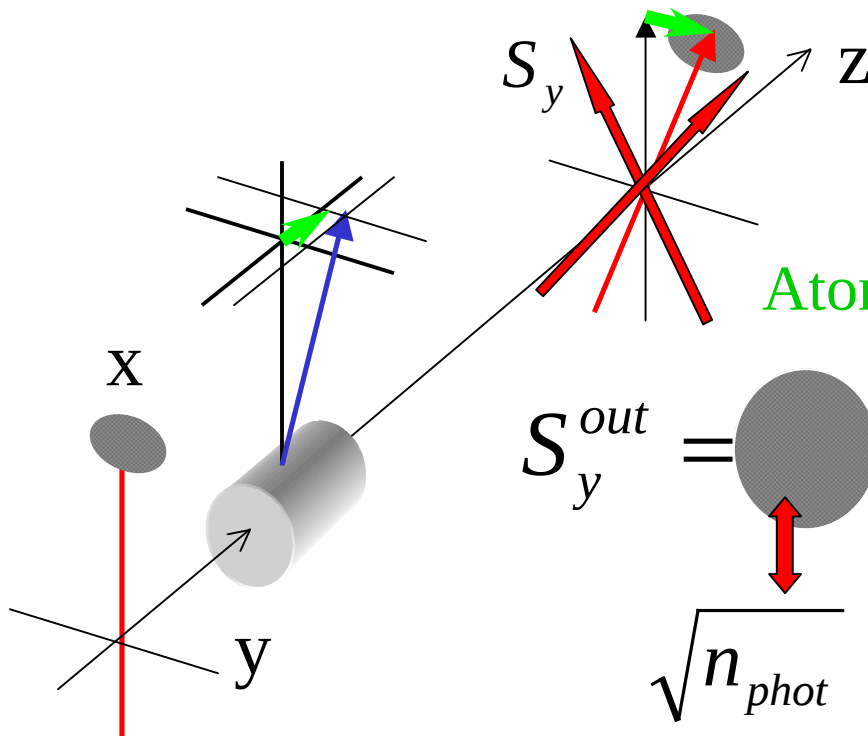
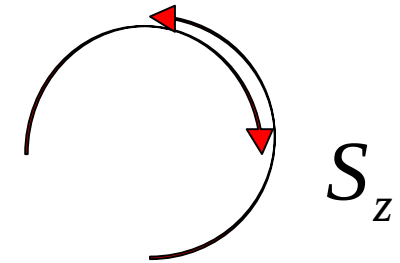
Off-resonant, perturbative:

$$\begin{aligned} H &\propto a_+^\dagger a_+ |g_-\rangle \langle g_-| + a_-^\dagger a_- |g_+\rangle \langle g_+| = \\ &= \frac{1}{2} (a_+^\dagger a_+ - a_-^\dagger a_-) (|g_-\rangle \langle g_-| - |g_+\rangle \langle g_+|) + \frac{1}{2} n_{phot} N_{atom} = \\ &= \frac{1}{2} S_z J_z + \frac{1}{2} \cancel{n_{phot} N_{atom}} \end{aligned}$$

Light / Atom - Interaction

Hamiltonian:

$$H = a S_z J_z$$



Faraday effect:

Atomic spins rotate polarization of light

$$S_y^{out} = \text{atom} + \kappa J_z \quad K=1$$

The diagram shows the output state S_y^{out} as the sum of the atom state and the coupling κJ_z . A red double-headed arrow labeled $K=1$ indicates the coupling strength.

$$\hat{X}_{light}^{out} = \hat{X}_{light}^{in} + \hat{P}_{atoms}^{in}$$

$$\kappa = \frac{1}{2} n_{phot} \frac{\sigma}{A} \frac{\gamma}{\Delta}$$

Figure of merit

$$\left(\frac{\kappa J_z}{S_y^{in}} \right)^2 \approx \alpha_{\Delta} s_{\Delta} \frac{\Delta^2}{\gamma^2} \gamma \tau_{pulse} = \alpha_0 s_{\Delta} \gamma \tau_{pulse} = \alpha_0 \eta$$

Probe depumping parameter:

$$\eta \ll 1$$

Light / Atom - Interaction

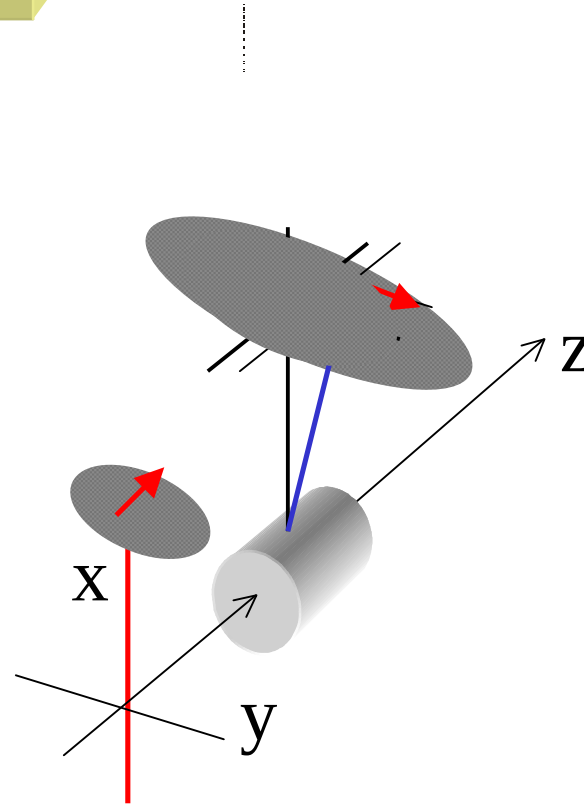
Back action of light on atoms:

Light rotates spins of atoms

$$J_y^{out} = J_y^{in} + kS_z$$

$$\updownarrow k=1$$

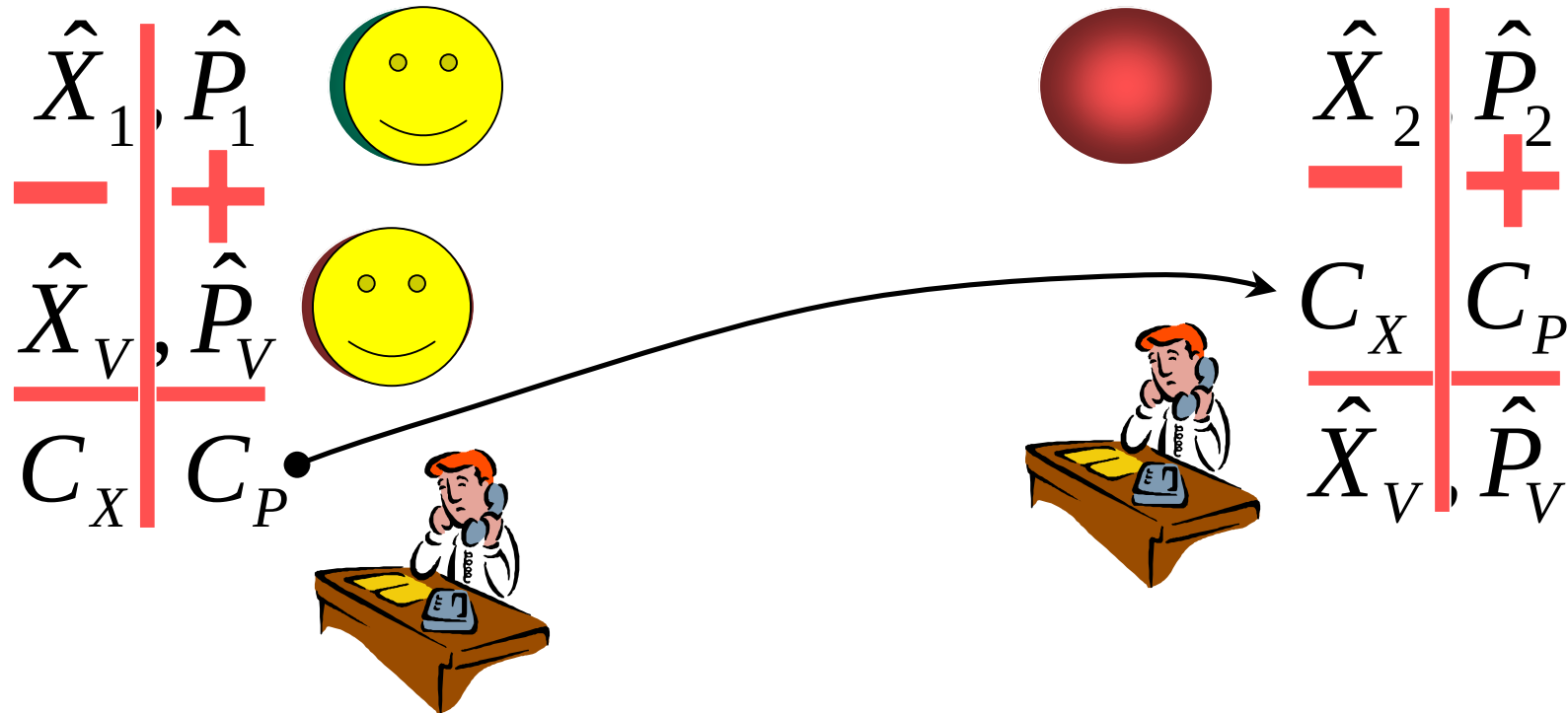
$$\hat{X}_{atoms}^{out} = \hat{X}_{atoms}^{in} + \hat{P}_{light}^{in}$$



Teleportation principle (continuous variables)

$$[X, P] = i, [X_1 - X_2, P_1 + P_2] = 0$$

L.Vaidman

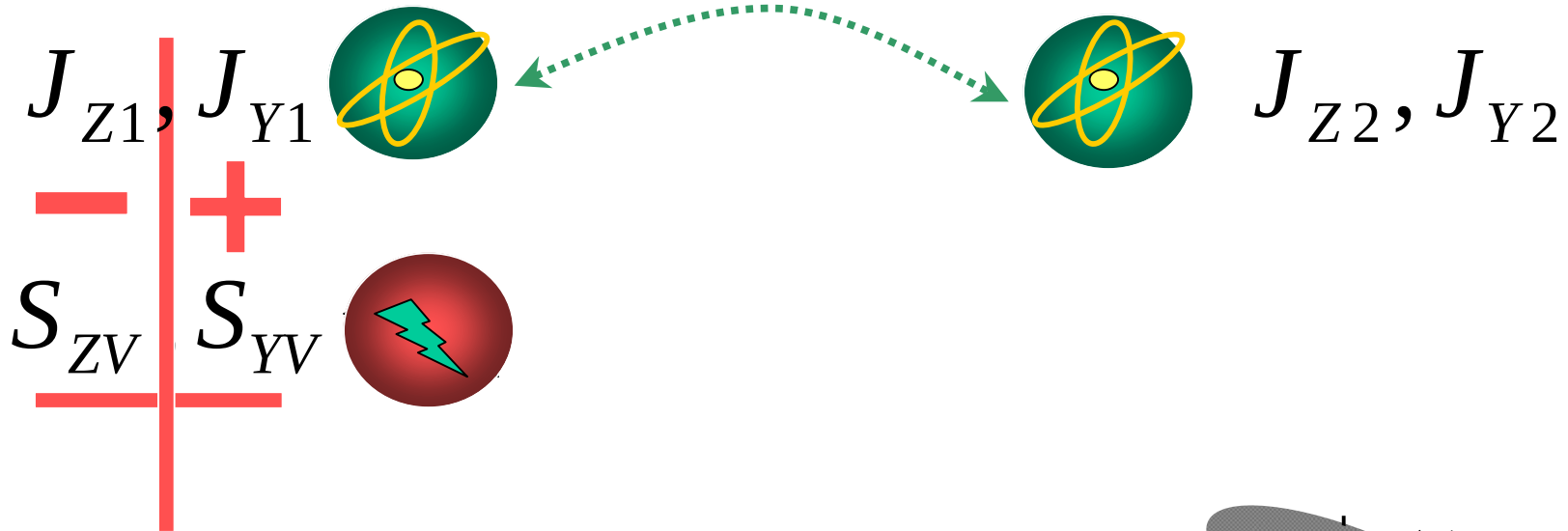


Demonstrated experimentally for light variables by Furusawa, Sørensen, Fuchs, Braunstein
Kimble, Polzik. *Science* 1998

Einstein-Podolsky-Rosen entangled state

$$X_1 - X_2 = 0, P_1 + P_2 = 0$$

Light-to-Atoms Teleportation

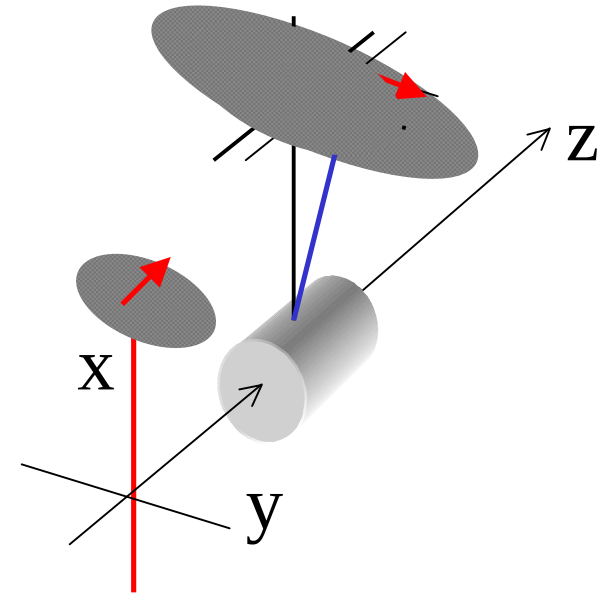


$$J_y^{out} = J_y^{in} + kS_z$$

$$S_y^{out} = S_y^{in} + kJ_z$$

$k=1$

$$\hat{X}_{light}^{out} = \hat{X}_{light}^{in} + \hat{P}_{atoms}^{in}$$

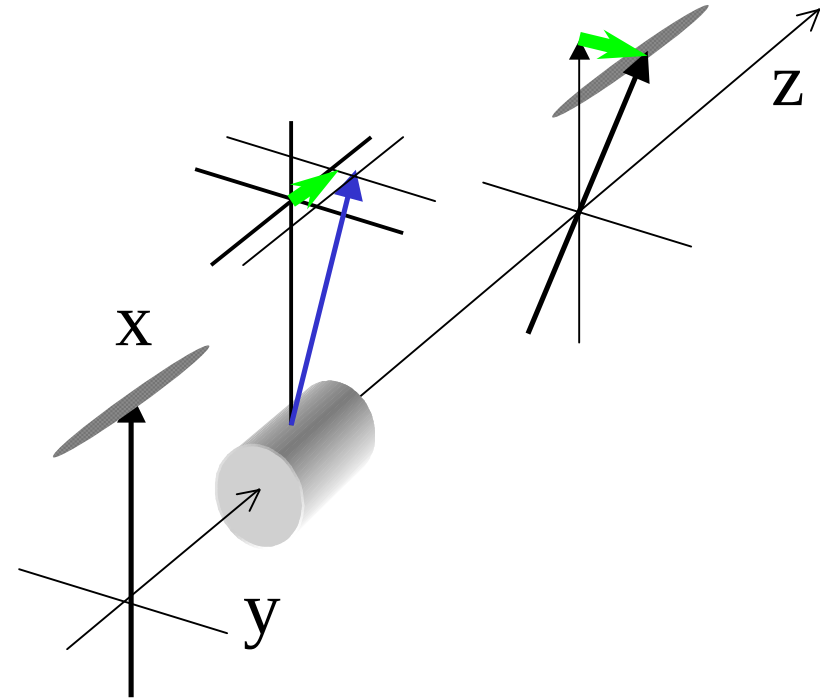


Quantum memory read-out: single pulse

Step 1

$$S_y^{out} = S_y^{in} + \frac{1}{2} an J_z = S_y^{in} + J_z^{in},$$

$$J_y^{out} = J_y^{in} + \frac{1}{2} a N S_z = J_y^{in} + S_z$$

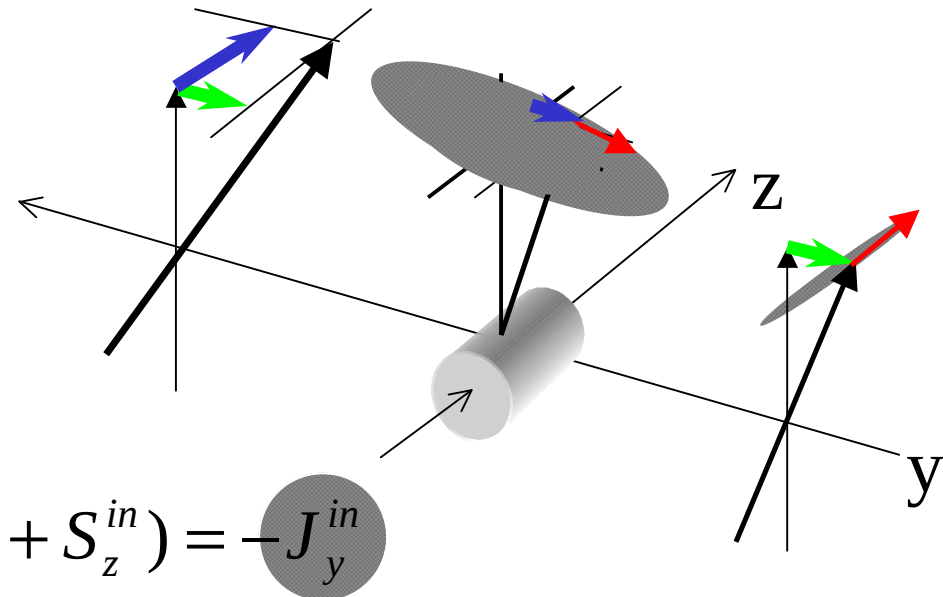


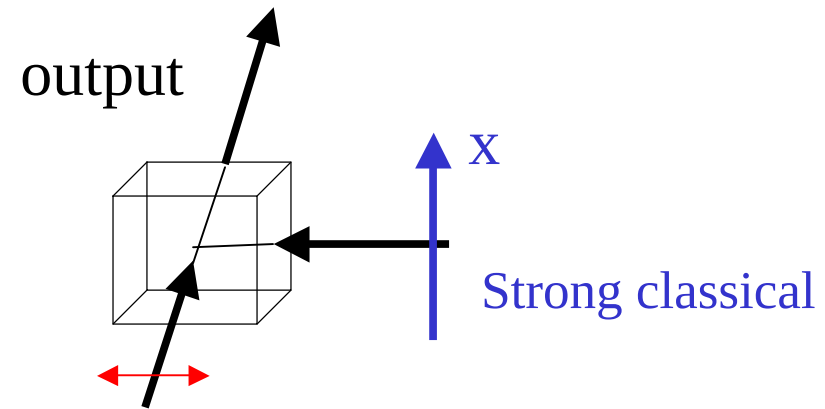
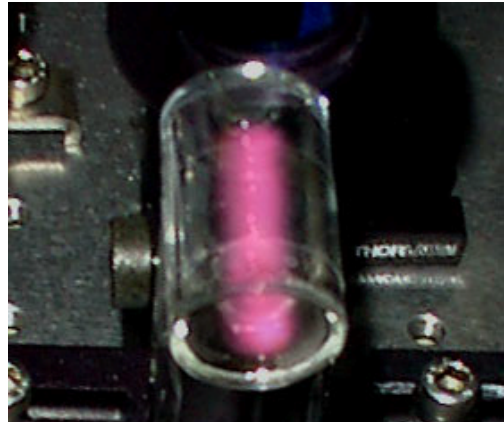
Step 2

Exchange y and z components:
pass light through $\lambda/4$ plate and
probe along spin-y axis

$$S_z^{out-2} = S_y^{in} + J_z^{in} \approx J_z^{in},$$

$$S_y^{out-2} = S_z^{out} - J_y^{out} = S_z^{in} - (J_y^{in} + S_z^{in}) = -J_y^{in}$$





Quantum:
Vacuum or squeezed

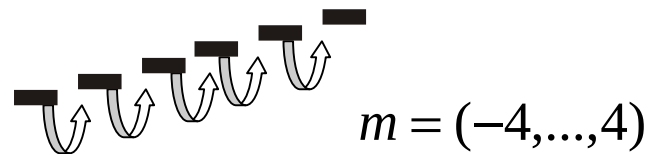
Light / Atom – Interaction, rotating frame

- To record a frequency component Ω of light we apply magnetic field $\mathbf{B}$ to atoms:

$$\dot{\hat{J}}_z(t) = \Omega \hat{J}_y(t) - \Gamma \hat{J}_z(t) + \hat{f}_z(t)$$

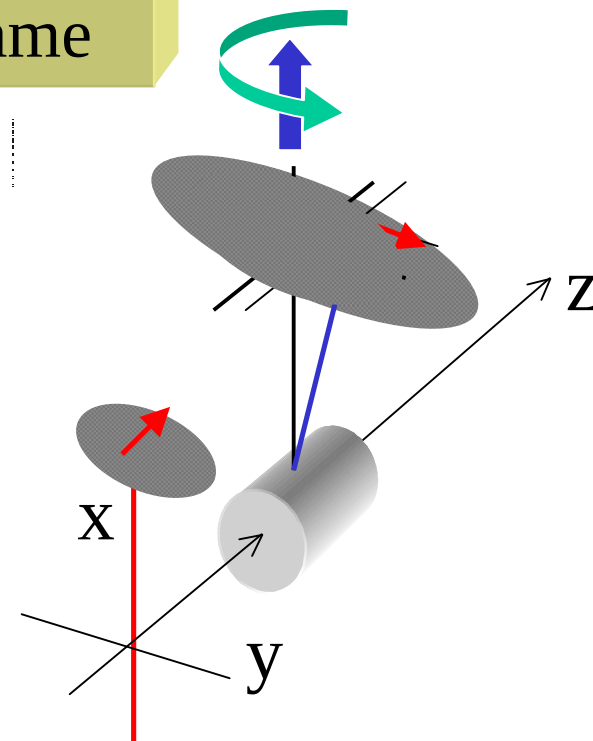
$$\dot{\hat{J}}_y(t) = -\Omega \hat{J}_z(t) - \Gamma \hat{J}_y(t) + k \bullet + \hat{f}_y(t)$$

$6S_{1/2}, F=4$

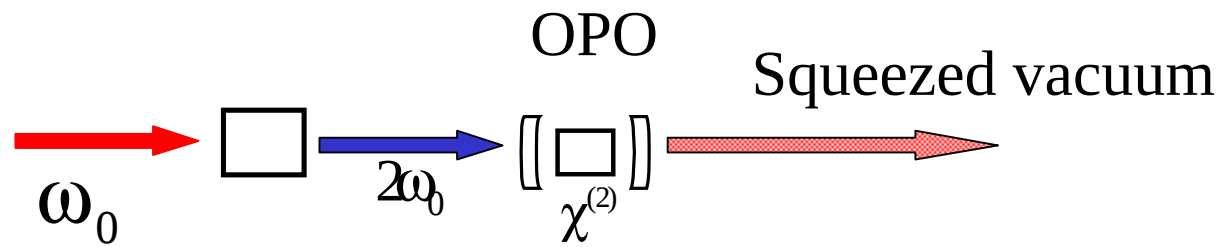


Measurement on light reads out atomic state which contains the memory term $\bullet$ centered around Zeeman frequency Ω

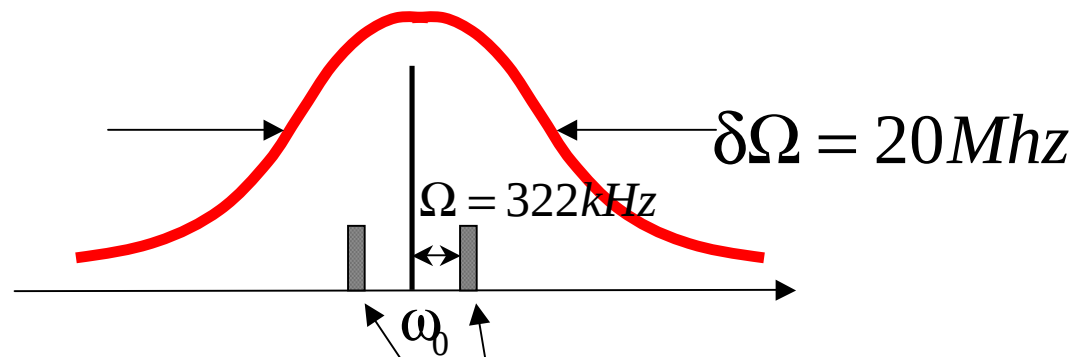
$$\hat{S}_y^{out}(t) = \hat{S}_y^{in}(t) + k \hat{J}_z(t)$$



Source of light

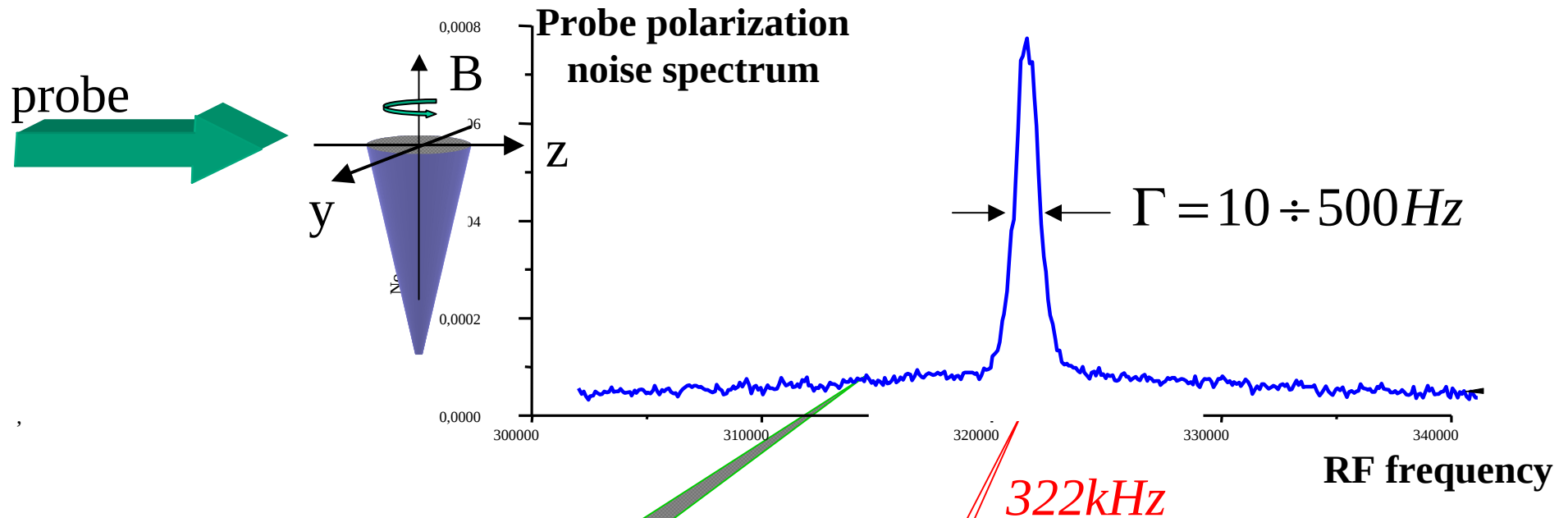


Spectrum



Sidebands recorded in atoms

Detecting quantum fluctuations of the spin

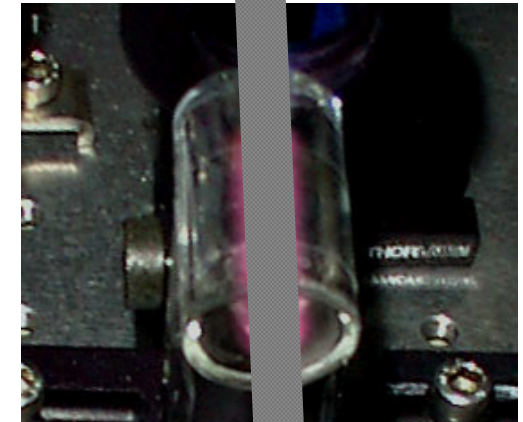
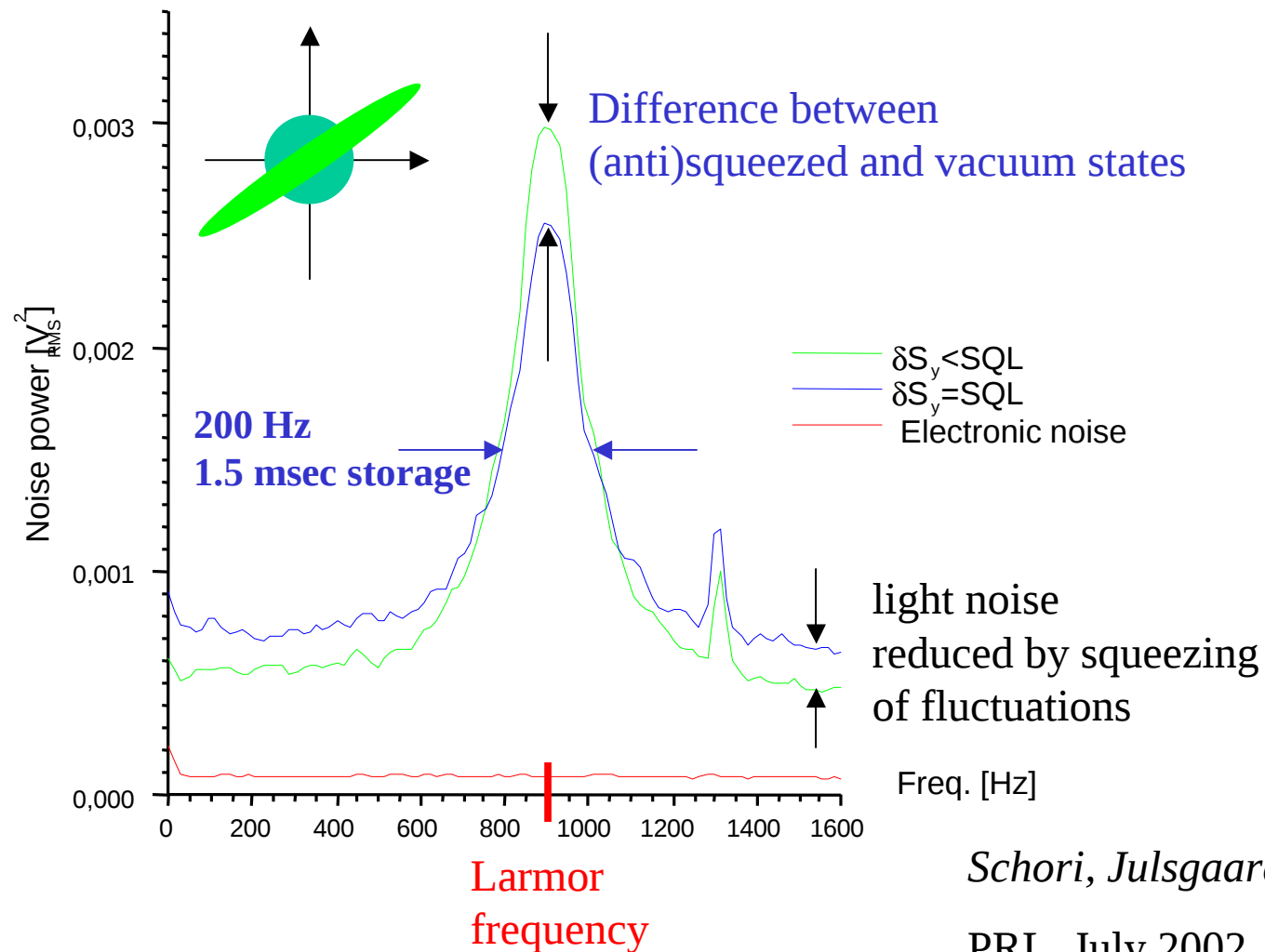


$$\hat{S}_y^{out}(\Omega) = X \hat{S}_y^{in}(\Omega) + \frac{a\Gamma}{\Gamma^2 + (\Omega - \Omega_L)^2} [\hat{J}_{z1} + bP \hat{S}_z^{in}(\Omega)]$$

$X \times P = 1$

(Partial) recording of quantum state of light in long-lived spins

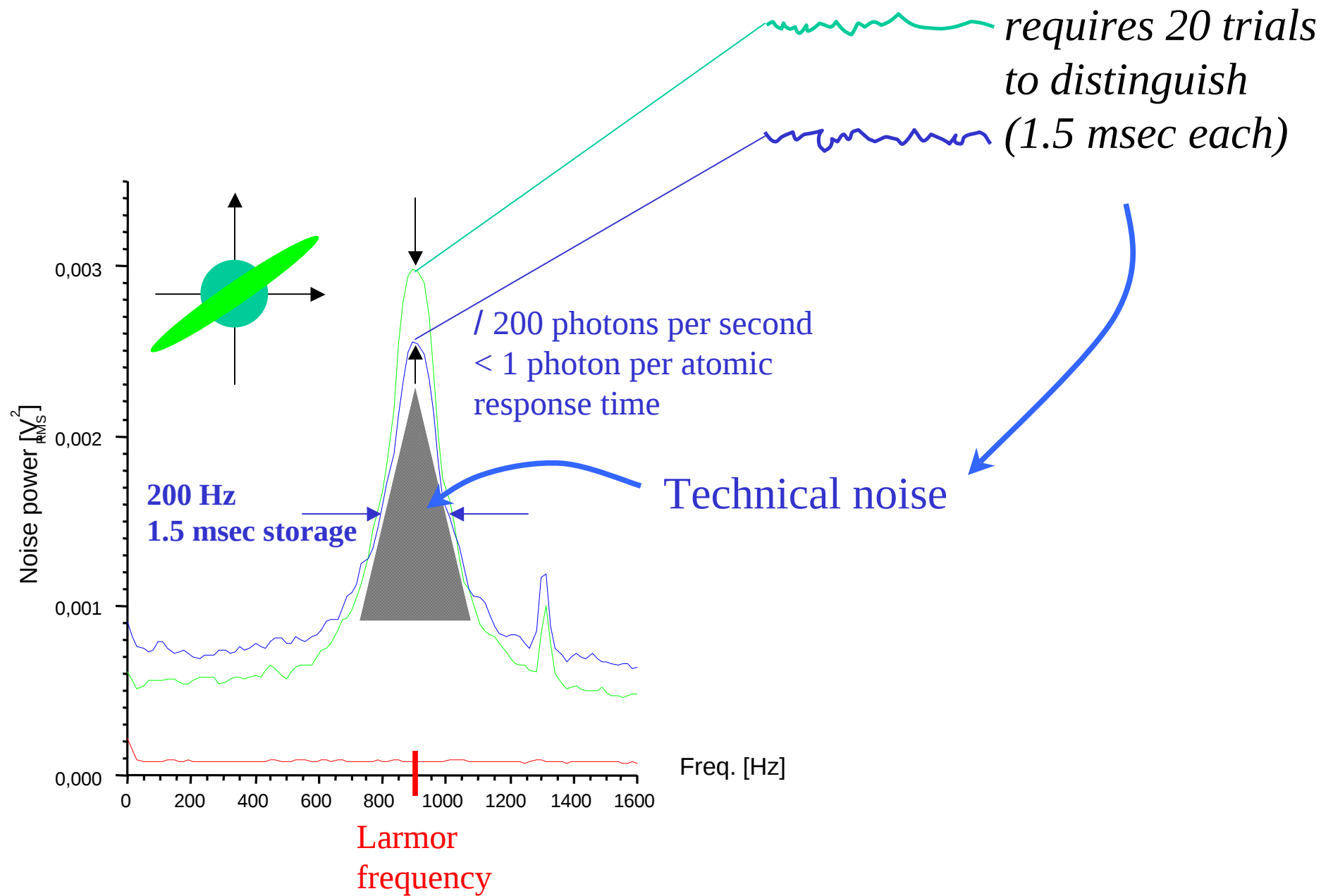
$$\hat{S}_y^{out}(\Omega) = X \hat{S}_y^{in}(\Omega) + \frac{a\Gamma}{\Gamma^2 + (\Omega - \Omega_L)^2} \left[\hat{J}_{z1} + bP \hat{S}_z^{in}(\Omega) \right]$$



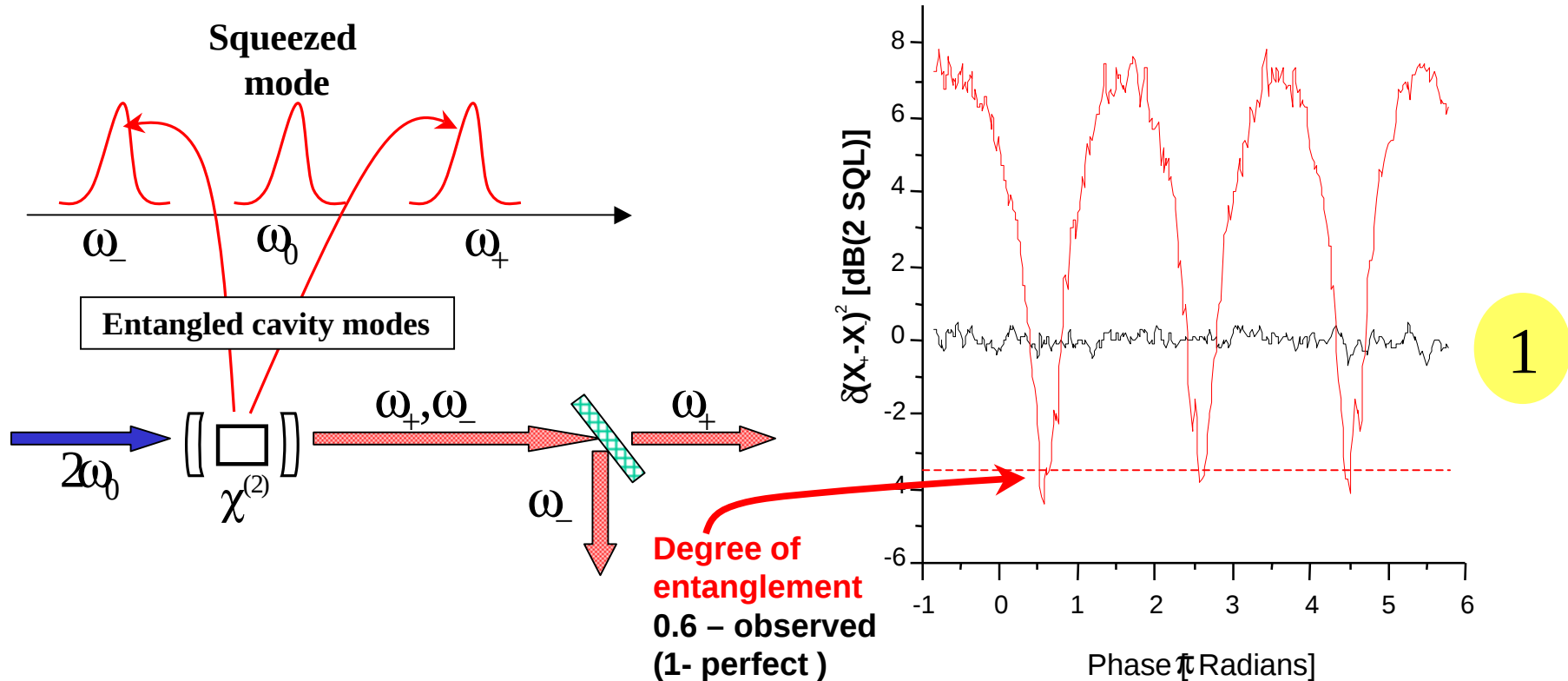
Light with
Controlled
X and P –
controlled
quantum state

Schori, Julsgaard, Sorensen, Polzik

PRL, July 2002, quant-ph/0203023



Narrowband tunable squeezed/entangled beams



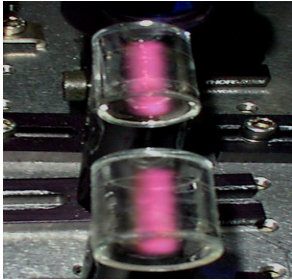
Sorensen, Schori, Polzik

To appear in PRA, quant-ph/0205015

$$\langle (X_1 - X_2)^2 \rangle + \langle (P_1 + P_2)^2 \rangle < 2$$

$$\text{OR } \begin{cases} \langle (P_1 + P_2)^2 \rangle < 1 \\ \langle (X_1 - X_2)^2 \rangle < 1 \end{cases}$$

Quantum communication protocols with entangled atomic samples and tunable EPR light



Entangled atomic samples
Nature 413, 400 (2001).



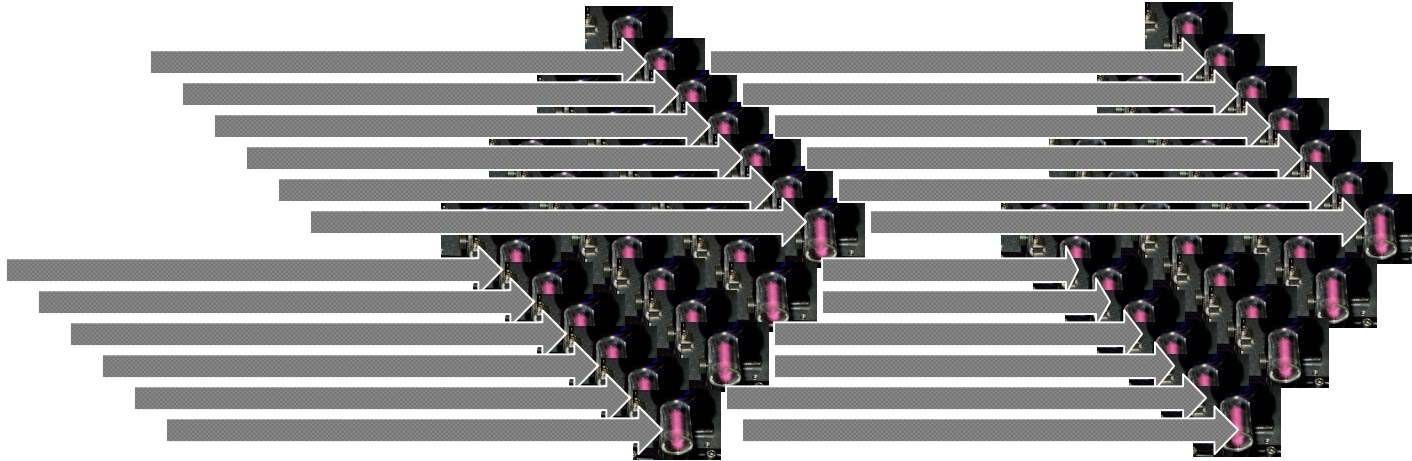
Entangled tunable light source
To appear in PRA, quant-ph/0205015

Protocols:

Quantum memory

Atomic teleportation

Scalability – an array of dipole traps or
solid state implementation – quantum holograms



Requirements:

- Long-lived 2 or more level state of a single particle
- The ability to optically pump the system into one of these states
- Dipole coupling to light
- Optically dense medium

Example: pencil-shaped cold cesium sample
50 microns by 500 microns
 10^5 atoms, 10^5 photons per pulse

Experiment: recording in the ground state coherence ($\tau=1.5\text{msec}$)

Schori, Julsgaard, Sorensen, Polzik

Phys. Rev. Lett. **89**, 057903 (2002),

quant-ph/0203023

Related work – entangled atomic ensembles

Julsgaard, Kozhekin, Polzik

Nature **413**, 400 (2001).

Theory - interspecies teleportation

Kuzmich, Polzik

Phys.Rev.Lett. **85**, 5639 (2000)

Earlier work: shot term recording in an excited atomic state ($\tau=32\text{ nsec}$)

Kuzmich, Mølmer, Polzik. *Phys.Rev.Lett.*, **79**, 4782 (1997)

Hald, Sorensen, Schori, Polzik. *Phys.Rev.Lett.*, **83**, 1319 (1999)